Dissipation in the turbulent ISM

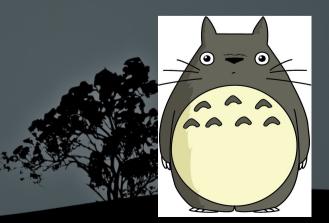
Edith Falgarone

François Boulanger, Benjamin Godard, Pierre Hily-Blant, François Levrier, Pierre Lesaffre, Guillaume Pineau des Forêts,

> Andrew Lehmann, Alba Vidal García, Thibaud Richard

> > **Totoro**





Molecules, magnetic fields and Intermittency in coSmic Turbulence Following the energy trail...

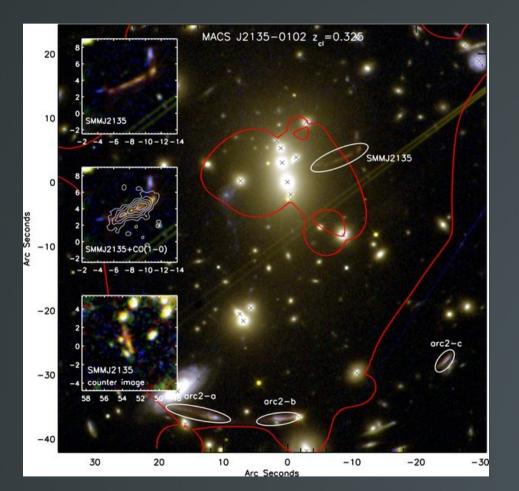


Bioluminescence in waves: plancton highlights Strong shear change

MISTy questions

- Origin of molecules in dilute and violent media ?
 - CO observed in diffuse irradiated media
 - Warm H2
 - extragalactic and galactic CH+
- Origin of the clumpy structure of the cold ISM ?
- Origin and Structure of the B field, its link with matter ?
- MHD turbulence dissipation (energy is dissipated in localised structures which affect the chemistry and magnetic fields)

The blink of the cosmic eyelash A galaxy seen through a gravitational lense



Cosmic Eyelash CO(1-0) Swinbank+2010 5 -5 -10 CH+(1-0) -15 -20 2 ՟ᢦ᠈ᢦᡁᠾᡄᢦ᠋᠓ᠸᠵ᠋ᢩ᠆ᢔᠮᢗᡆ᠋ᠺᡀᢦᡀᢧᠧᡔᢍᡟ_{ᡡᡗᡗ}ᢍᡙᡗᢦ᠆ -2000 -1000 1000 2000 velocity (km s⁻¹)

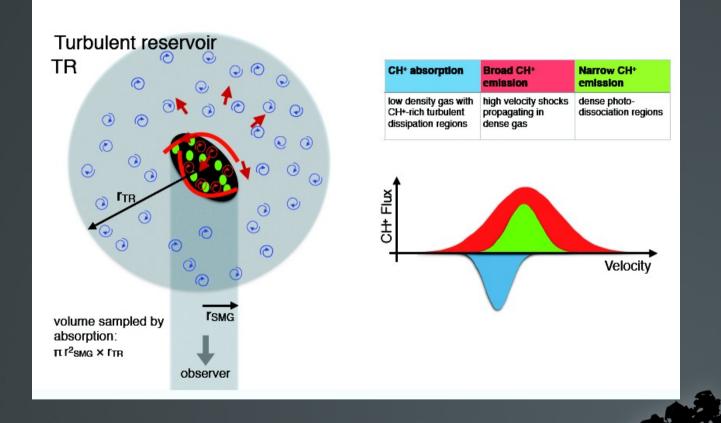
Swinbank+2010 Redshift z=2.33



Observational Context

- Broad lines of CH+ in emission and absorption in high-z galaxies (Falgarone+17 Nature, 548, 430)
- Molecular emission in colliding galaxies (Stefan's Quintet)
- Molecular emission hints at very dense and cool, probably clumpy media
- Broad lines: equipartition between large scale thermal energy (10⁷ K) and kinetic energy of dense cool gas.

The blink of the cosmic eyelash

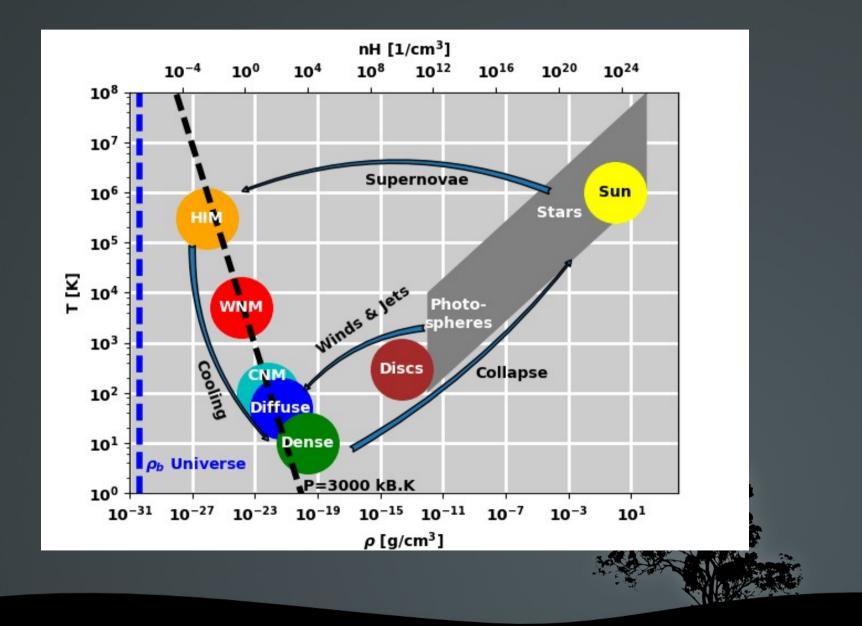


Falgarone+2017

Interpretation in MIST The energy trail(s)

- Large scale energy is transferred to smaller scales by turbulent cascade (stretching and compression) and by cooling cascade (condensation)
- The cooling cascade proceeds with a phase separation. How is kinetic energy shared between phases ?
- Energy is radiated in bursty dissipative structures
- Molecules are produced and excited by these dissipative structures (e.g. vortex, shocks)

The phases of the ISM



Typical values

- Huge dynamical range of length scales,
- but not so big for velocity

1/

Reynolds number: UxL / dissipation coefficient

		HIM	WNM	CNM	Diffuse	Dense	Discs	Sun
_	Density ρ [cm ⁻³]	0.004	0.6	30	200	10^{4}	10^{10}	1 g.cm^{-3}
	Temperature T [K]	3.10^{5}	5000	100	50	10	300	10^{6}
	Length scale L [pc]	100	50	10	3	0.1	200 AU	5.10^{-3} AU
	Velocity U [km.s ⁻¹]	10	10	10	3	0.1	0.1	1
	\mathcal{M}	0.2	2	13	7	0.5	0.1	0.02
	$\mathcal{M}_\mathcal{G}$	130	20	15	6	0.8	0.08	0.003
	${\mathcal R}$	10^{2}	10^{5}	10^{7}	10^{7}	10^{6}	10^{9}	10^{17}
	\mathcal{R}_m	10^{21}	10^{20}	10^{18}	10^{17}	10^{15}	10^{9}	10^{10}
	\mathcal{R}_{AD}	10^{3}	10^{3}	10^{2}	10^{3}	10^{4}	10^{5}	10^{20}
	Ionisation fraction	1	10^{-2}	10^{-4}	10^{-4}	10^{-4}	10^{-7}	1
$M_{10x}^{Mass per ion [amu]}$ $N_e B_z [10^{19} \text{cm}^{-2} 3 \mu \text{G}]$		1	1	12	12	12	24	1
./ 工 ($N_e B_z \ [10^{19} \text{cm}^{-2} 3 \mu \text{G}]$	0.12	0.09	0.009	0.01	0.02	0.2	3×10^{15}
1/10x $N_e [10^{19} \mathrm{cm}^{-2}]$		0.05	0.5	0.3	0.3	0.1	10^{3}	10^{27}

Dissipation in decaying turbulence (incompressible runs)

 $n_{H} \sim 100/cm^{3}$ $<u^{2}>\sim<b^{2}/\rho>$ $Re=LU/v \sim 2.10^{7} 10^{3}$ $Re_{m}=LU/\eta \sim 2.10^{17} 10^{3}$ $Re_{AD}=L/U/t_{AD} \sim 10^{2}$

1 pc

Line of sight integrated dissipation:

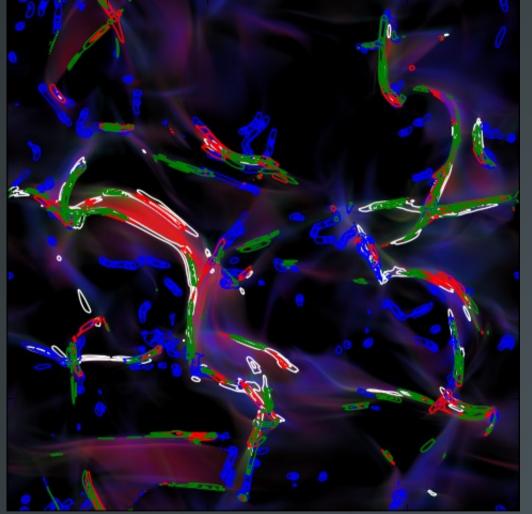
$$\epsilon_{\rm diss} = \nu \rho S_{ij}[u] \partial_i u_j + \eta |\nabla \times B|^2 + |F_{in}|u - v|^2$$

(Momferratos PhD thesis: 512³ spec. elts Incompressible simulations by ANK, pseudo-spectral code with AD)

Giorgos' PhD: Decaying Incompressible MHD + AD

- Dissipation localised on sheets, structure extraction
- Measured statistics of dissipative structures (PDFs and correlation between characteristic prop.^{ties})
- Each sheet is 'pure' in its dissipation nature: viscous, ohmic or ambipolar heating
- Correlations between dissipation proj. and increments
- Initial conditions matter a lot
- B field orientation is not random w.r.t. to dissipation in slices
- A.D. Forces jxb=0 from small scales towards large

Observable increments vs. dissipation Lbox/2



• <u>Background</u>: Dissipation rates Ohmic Viscous AD

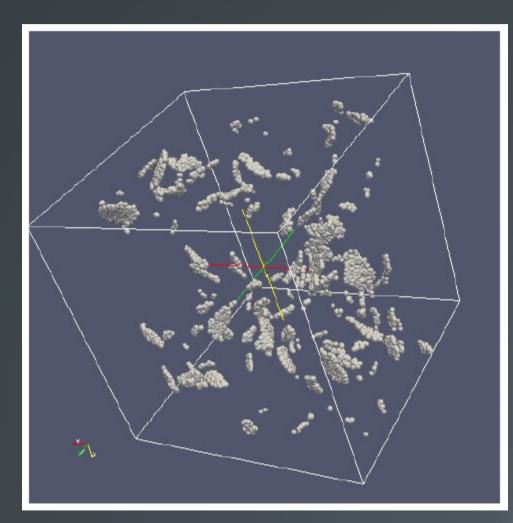
• <u>Contours:</u> Increments of

integrated observables:

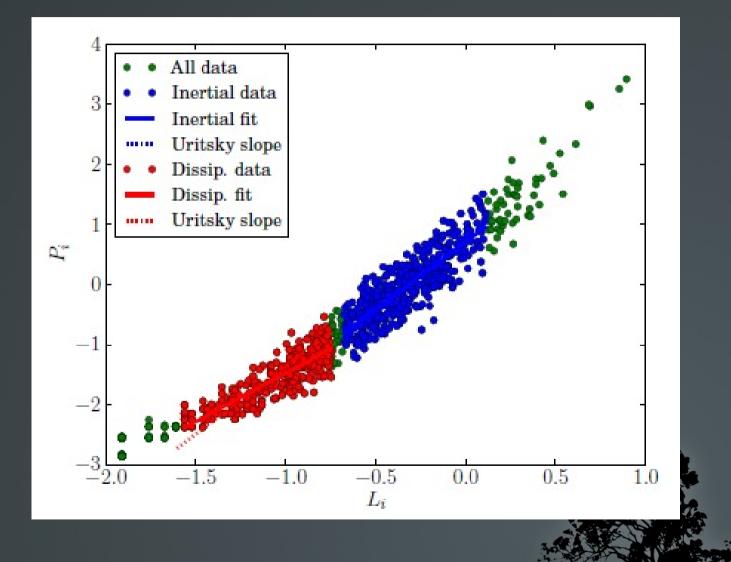
- LOS velocity (white)
- Stokes Q (green)
- Stokes U (red)
- POS polarisation angle (blue)

<u>NOTE:</u> increment of polarisation angle (blue contours) are less correlated to dissipation. Better use Q,U.

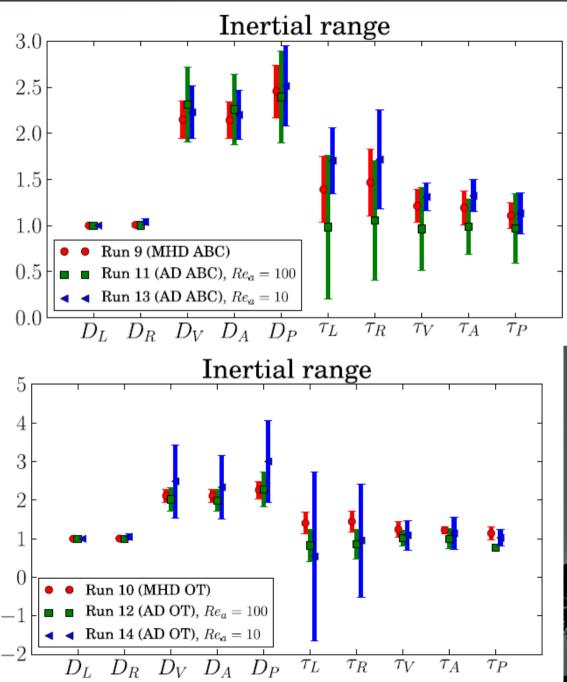
Dissipative structures (locii of intense dissipation, > μ +2 σ)



Structure statistics (one example)

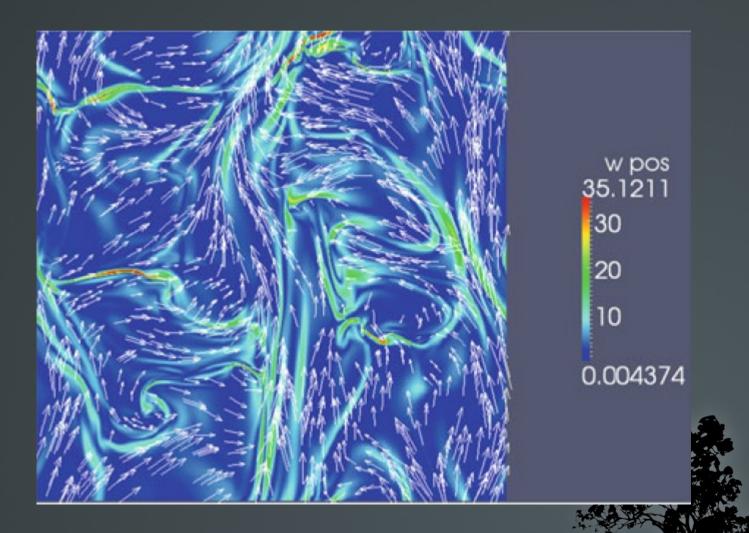


Initial conditions matter





Slice of B(p.o.s.) and |curl(u)| relative orientation not random



JxB ~ 0 with ambipolar diffusion

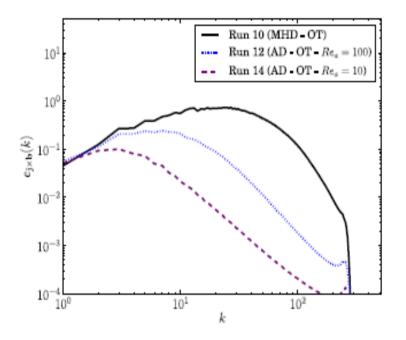


Figure 4.6: Power spectra of $\mathbf{j} \times \mathbf{b}$ for high resolution runs 10,12 and 14. The field becomes force-free at small scales when the strength of the AD is increased.

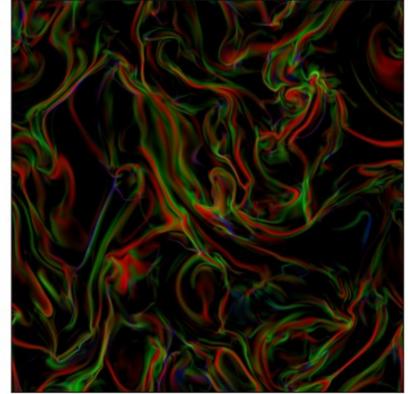
Ques: does this mean that at small scales, Hydrodynamic turbulence prevails ? And Bfield is force-free (jxb=0), independent ? Maybe not: induction equation must be compatible with force-free.

Simulations of decaying turbulence. Compressible: Isothermal 3D MHD (Mach 4, ABC)

 $n_{H} \sim 100/cm^{3}$ $<u^{2}>\sim<b^{2}/\rho>$ Re=LU/ $\nu \sim 2.10^{7} \ 10^{3}$ Re_m=LU/ $\eta \sim 2.10^{17} \ 10^{3}$ (1020³ pixels)

Heating nature in decaying MHD turbulence

~1 pc



Red: Ohmic heating Blue: 4/3 v div(u)² Green: v curl(u)²

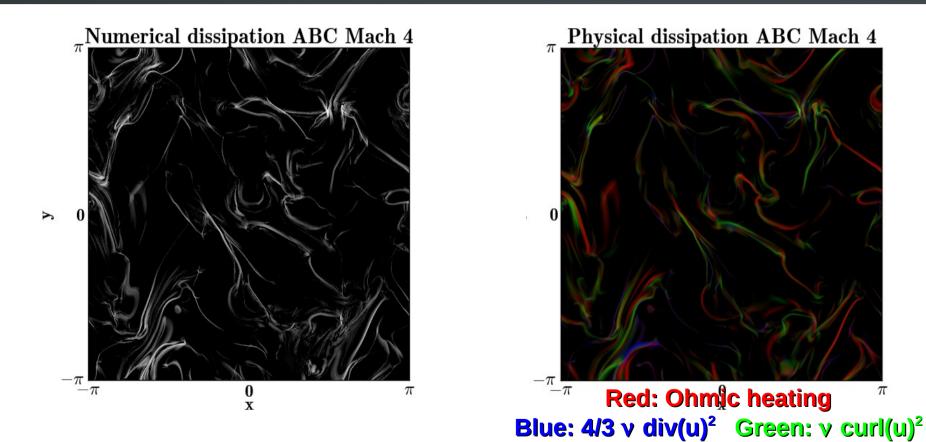
(Momferratos PhD thesis: DUMSES simulations with careful treatment of viscous and resistive dissipation)

Thibaud's Internship (& PhD) Decay Compressible MHD (no AD)

- Nature of dissipation is mixed
- [grid => need to recover dissipation from grid]
- B field // dissipative structures in 3D (but not in projection)
- Dissipation is dominated by low convergence



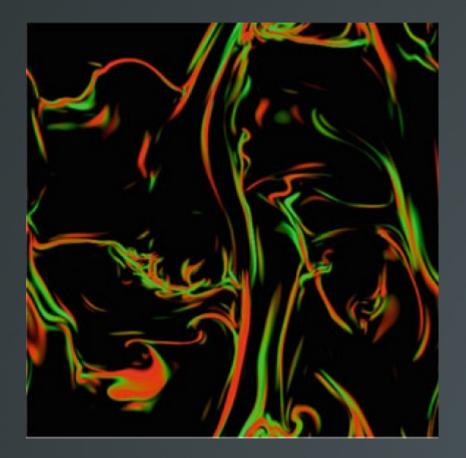
Grid Dissipation



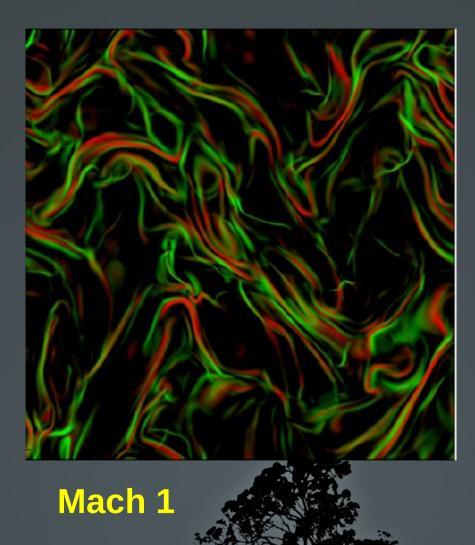
Need to estimate *numerical* dissipation to recover the properties of the *total* dissipation



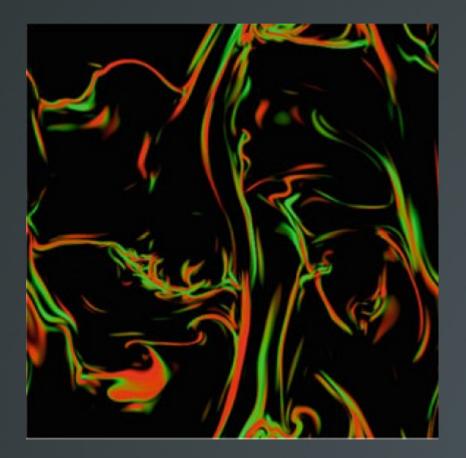
Now ohmic and viscous dissipation are mixed



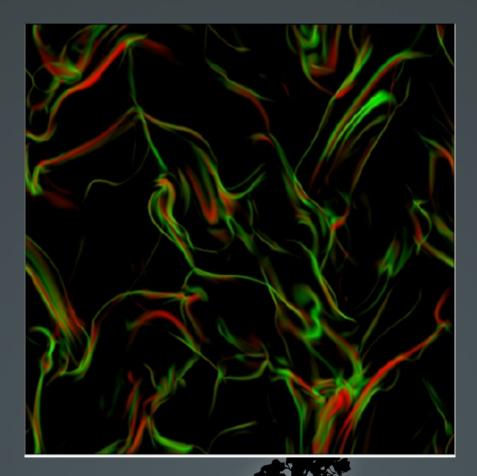
Mach 0



Now ohmic and viscous dissipation are mixed

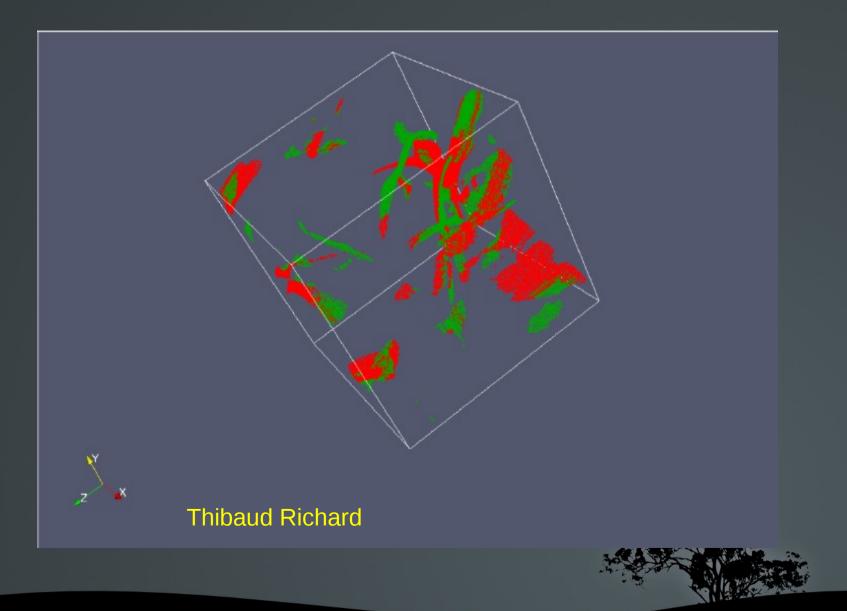


Mach 0



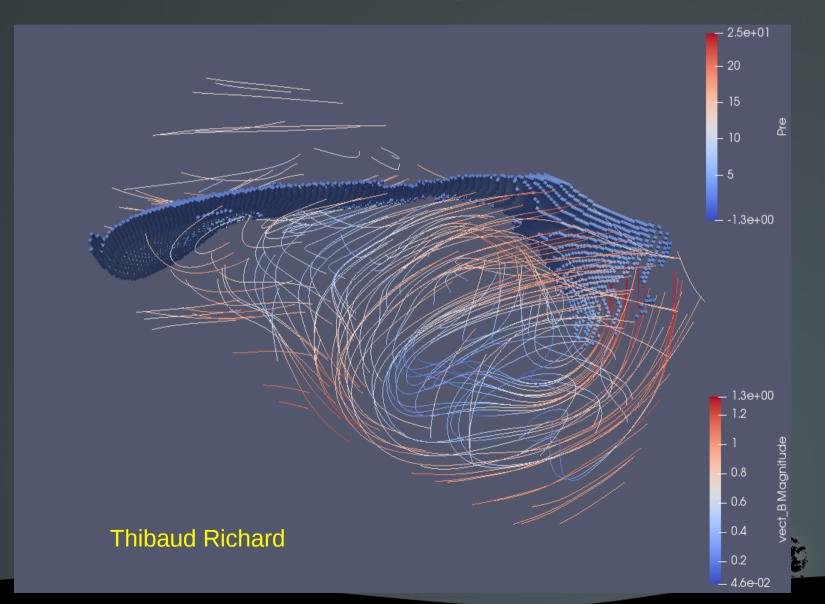
Mach 4

Viscous and **Ohmic** dissipation



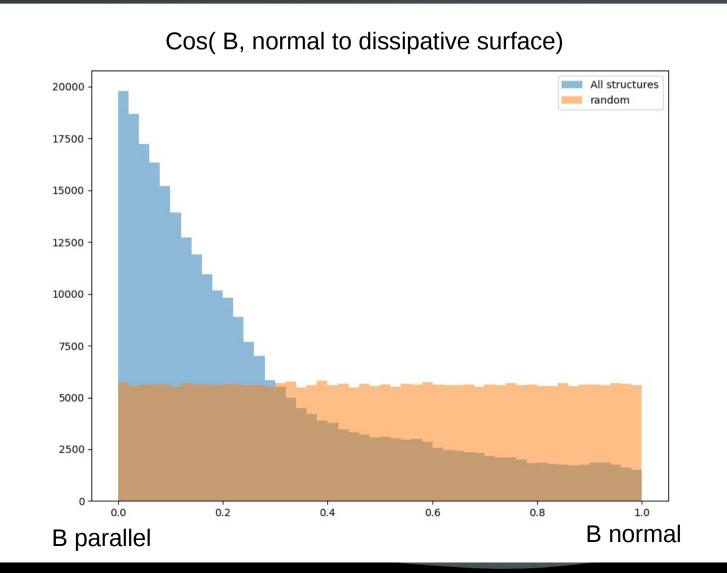
Dissipative structures extraction

Find connected sets where dissipation > mean + 2.std



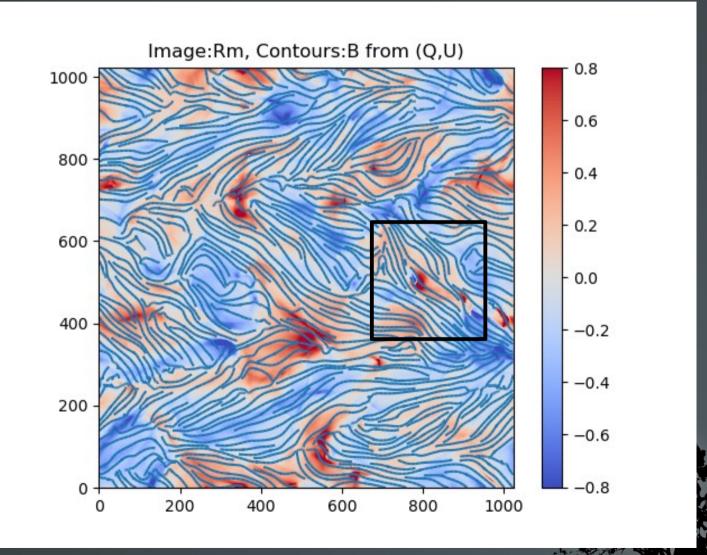
B field is mostly parallel to structures

Thibaud Richard



Projection over simulations

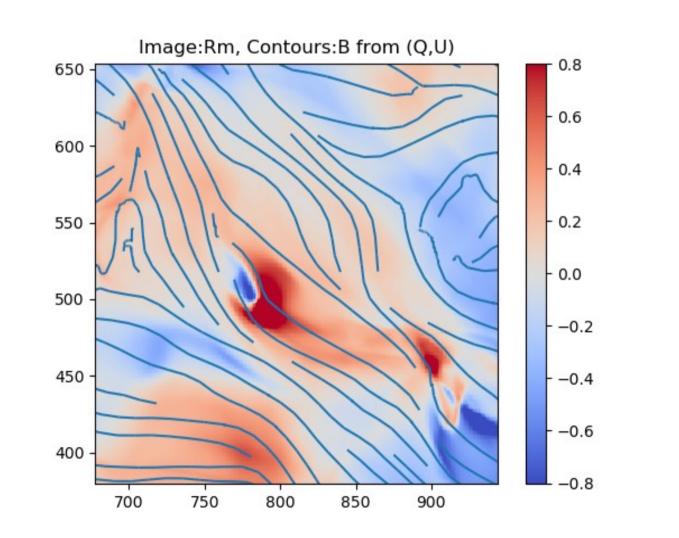
Rotation measure overlayed with p.o.s. B field direction



Mach 4, ABC

Projection over simulations

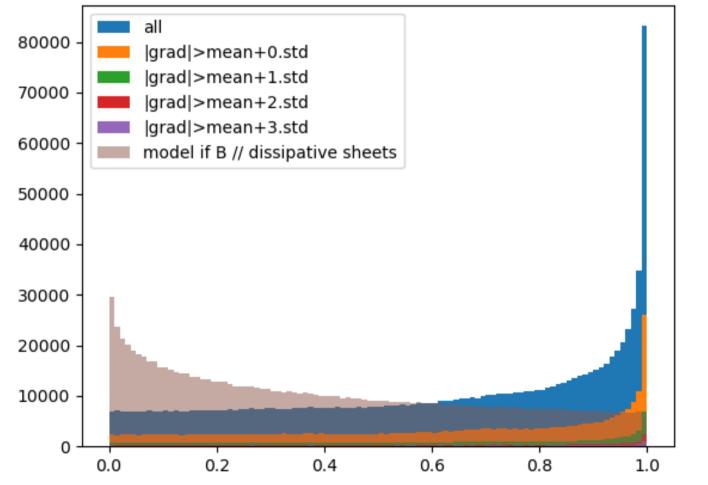
Rotation measure overlayed with p.o.s. B field direction



Mach 4, ABC

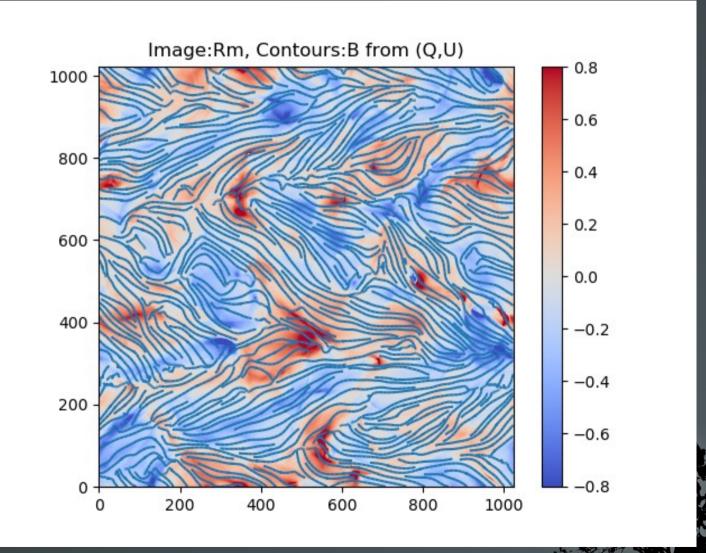
Orientation statistics

cos(angle) between grad(Rm) and B(Q,U)



Projection over simulations

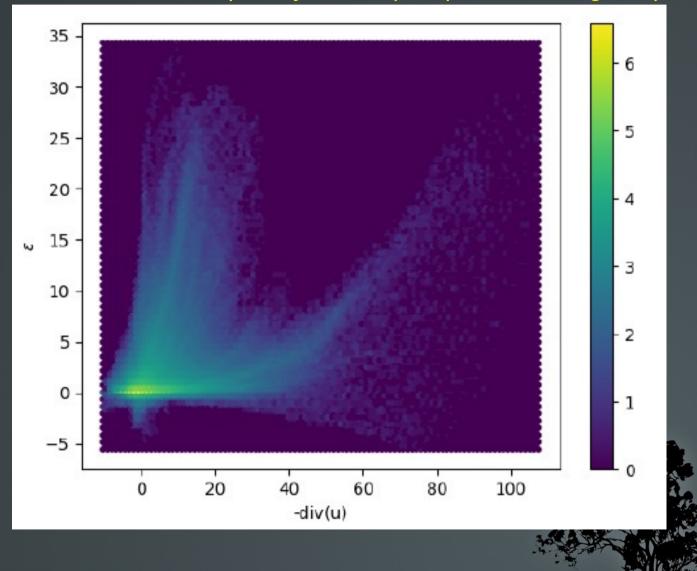
Rotation measure overlayed with p.o.s. B field direction



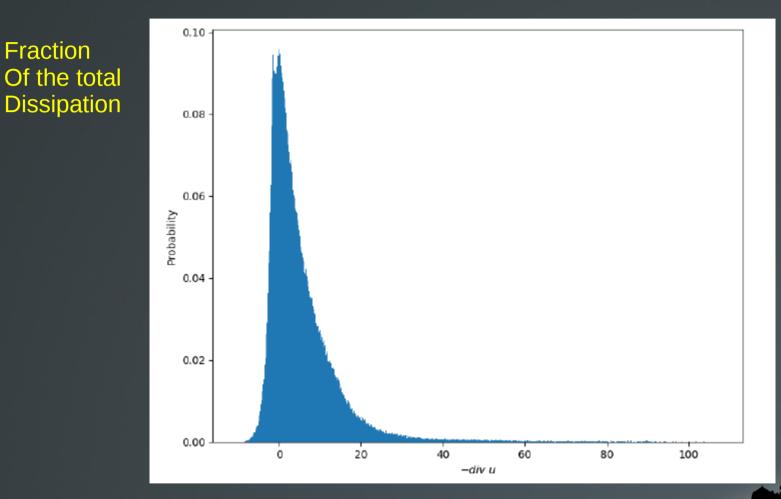
Mach 4, ABC

Dissipation is mostly incompressible

Volume fraction occupied by bins of (dissipation.convergence)



Dissipation is mostly incompressible



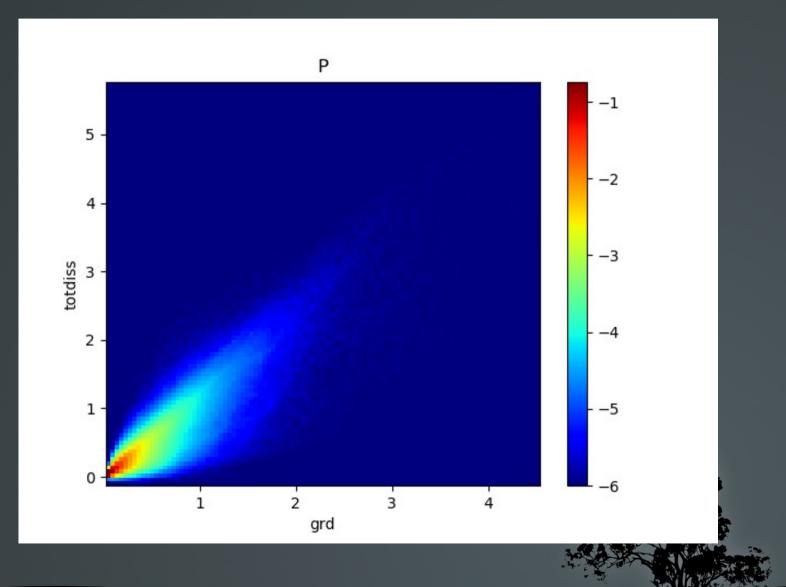
Convergence (unitless)

Pierre (+Thibaud, Andrew, Ben)

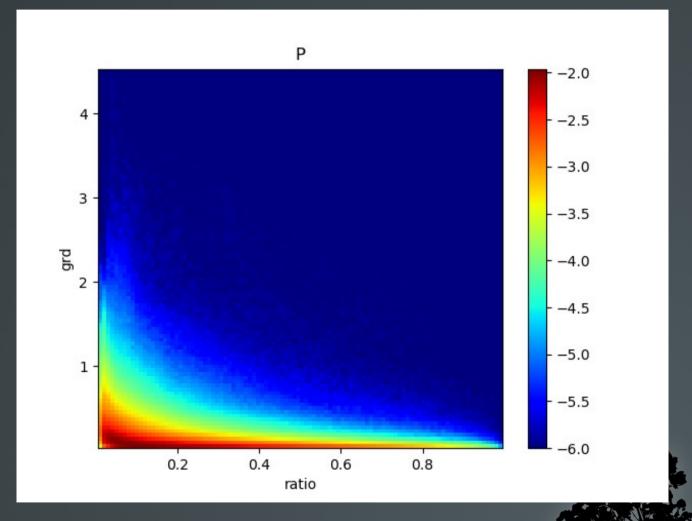
- Look at MHD variables: w=(rho, u, b) and search for direction of max gradient of w.
- Large gradients have a well determined orientation (meaning: plane-parallel is OK)
- Decompose gradients in fast/intermerdiate/slow waves
- Result: OK for single planar shocks, messy in compressible turbulence...

TODO: get shocks and parameters, measure stats, shock collisions ?

Total Dissipation vs Gradient



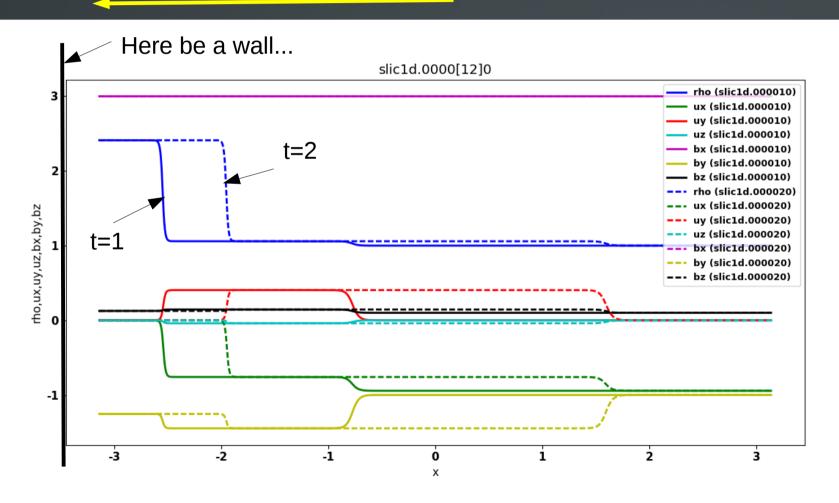
Isotropy (a2/a3) vs Gradient



Test Slow & Fast shock

Gas flows from right onto a wall on the left. 2 snapshots shown.

Direction of the inflow



Isothermal MHD equations

$$\frac{\partial \rho}{\partial t} + \boldsymbol{\nabla} \boldsymbol{.} (\rho \boldsymbol{u}) = 0$$

$$\partial_t \rho u_i + \partial_j (\rho u_i u_j + p \delta_{ij}) - \boldsymbol{j} \times \boldsymbol{b}_i + \partial_j (\rho \nu S_{ij}[u]) = 0$$

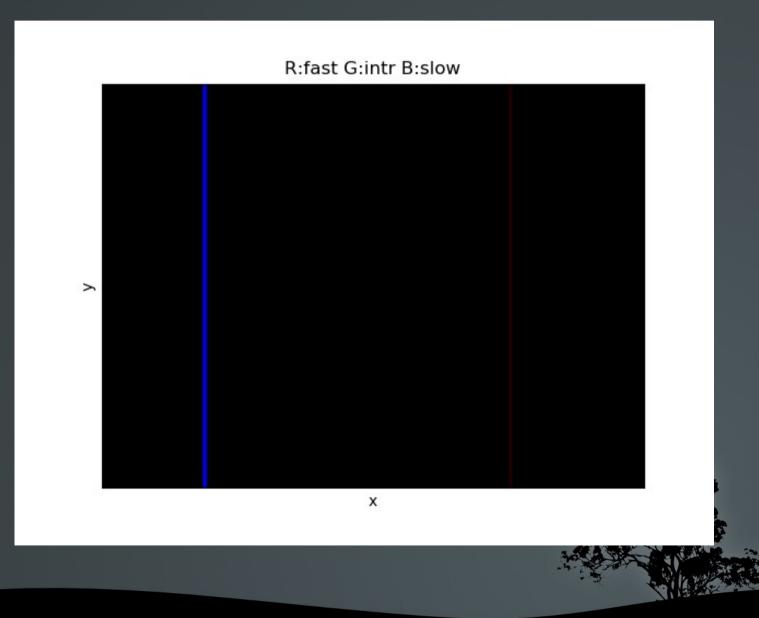
 $\partial_t \boldsymbol{b} - \boldsymbol{\nabla} \times (\boldsymbol{u} \times \boldsymbol{b}) + \boldsymbol{\nabla} \times (\eta \boldsymbol{j}) = \boldsymbol{0}$

$$j = oldsymbol{
abla} imes b$$

$$p = \rho c^2$$

$$S_{ij}[u] = \frac{1}{2}(\partial_i u_j + \partial_j u_i) - \frac{1}{3}\partial_k u_k \delta_{ij}$$

Gradient Decomposition

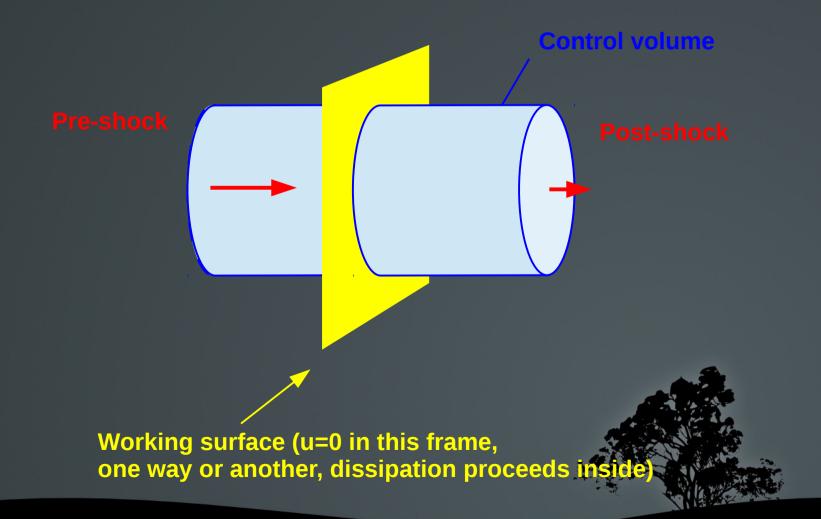


Gradient Decomposition in MHD compressible turbulence



Rankine Hugoniot

Flux conservation through a *steady* planar shock



Rankine Hugoniot

 Conservation of mass, momentum and magnetic flux in the steady shock *frame* induces relationships between preshock and post-shock physical conditions.

$$\begin{bmatrix} B_x \end{bmatrix}_{pre}^{post} = 0$$
$$\begin{bmatrix} \rho \ u_x \end{bmatrix}_{pre}^{post} = 0$$
$$\begin{bmatrix} (B \times u)_y \end{bmatrix}_{pre}^{post} = 0$$
$$\begin{bmatrix} (B \times u)_z \end{bmatrix}_{pre}^{post} = 0$$
$$\begin{bmatrix} \rho u_x u_y - B_x B_y \end{bmatrix}_{pre}^{post} = 0$$
$$\begin{bmatrix} \rho u_x u_z - B_x B_z \end{bmatrix}_{pre}^{post} = 0$$
$$\begin{bmatrix} \rho u_x^2 + P \end{bmatrix}_{pre}^{post} = 0$$

• Examples:

* Compression = Mach² in an isothermal shock * Max temperature $\sim u^2$ expresses conversion of kinetic to thermal energy in a viscous front $T_{\rm max} = 53 \,{\rm K} \,(u/1 \,{\rm km \ s^{-1}})^2$

For the molecular weight of the ISM:

Questions, ToDo

- What can we do to link to observations ?
- Single planar shock ? Single curved shock (Tram)?
- Link to Faraday rotation ?
- Collection of shocks ↔ Chemistry ?
- Focus of one structure in 3D MHD: how to reconstruct emissivity ?
- Test SHOCK_FIND. Improve on it ?
- Statistic theory of shocks ?
- Contact Discontinuity detection ↔ condensation fronts in thermal instability simulations

TODO

- Color 2d histo with div (u) / curl(u).
- Back to obs values for div(u), eps.
- Oblique shocks in Paris-Durham
- Statistics of slow / fast , convolve with Paris-Durham code => JWST predictions.
- Erwan: beware of scales of gradient: choose relevant ones for energy ?

